

A Short Rugged Ferrite Half-Wave Plate for a Single-Sideband Modulator*

One type of microwave phase shifter consists of two quarter-wave plates between which is placed a half-wave plate whose principal axis is rotatable.¹ The phase shift introduced is directly proportional to the angular displacement of the principal axis of the half-wave plate. Continuous rotation of the principal axis causes continuous advancement or retardation of the phase of the signal traversing the phase shifter. This causes a frequency shift of the signal and the device can thus be used as a single-sideband modulator.²

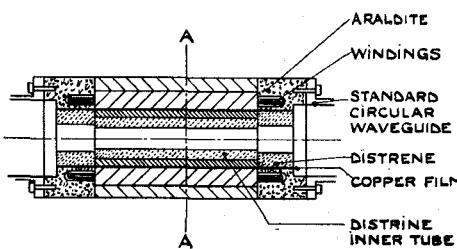


Fig. 1—Cross section of ferrite half-wave plate.

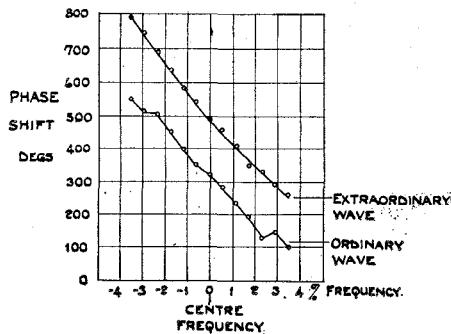


Fig. 2—Phase shift vs frequency characteristic or ordinary and extraordinary wave propagation.

A ferrite half-wave plate in reduced guide has been described by Karayianis and Cacheris.³ The direction of the applied transverse magnetic field forms the principal axis of the half-wave plate. The advantage of working in reduced guide is that since the guide is more dispersive, a higher differential phase shift per applied field can be obtained. The major disadvantage is the increased difficulties associated with matching the ferrite loaded reduced guide to the normal waveguide run. Matching was attempted by Cacheris and Karayianis by using two 2-inch dielectric tapers. These are difficult to manufacture and are lengthy for some applications.

A certain amount of work has been carried out using a tube of ferrite in reduced guide. The ferrite used was an experimental ferrite Type MM3 supplied by Marconi's Wireless Telegraph Company, Ltd. The dimensions of the tube are 0.7 inch O.D. \times 0.5 inch I.D. \times 2 inches long. The ferrite tube was loaded with distrene tubes of 0.5 inch O.D. and various I.D. dimensions. The final half-wave plate arrangement is shown in Fig. 1.

Fig. 2 shows a typical phase shift characteristic obtained. The kinks associated with the ordinary wave phase characteristic are a function of the ferrite tube dielectric loading and the input matching arrangement. The input circuit consisted of a single

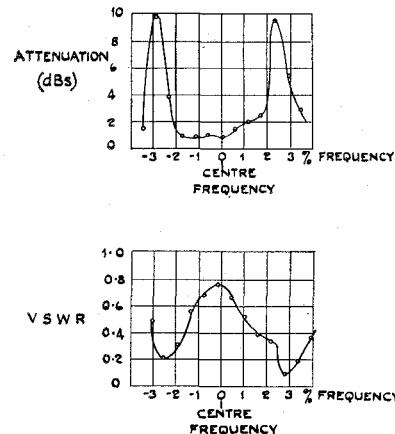
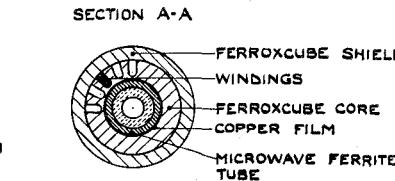


Fig. 3—Attenuation and VSWR vs frequency characteristic of the ferrite half wave plate.

* Received by the PGMTT, October 9, 1958.
¹ A. G. Fox, "An adjustable wave-guide phase changer," *Proc. IRE*, vol. 35, pp. 1489-1498; December, 1947.

² J. Cacheris, "Microwave single-sideband modulator using ferrites," *Proc. IRE*, vol. 42, pp. 1242-1247; August, 1954.

³ N. Karayianis and J. Cacheris, "Birefringence of fibers in circular waveguide," *Proc. IRE*, vol. 44, pp. 1414-1421; October, 1956.

The author wishes to thank the manager, The English Electric Company, Ltd., Luton, for permission to publish this note and Dr. Benzie of Marconi's Wireless Telegraph Company, Ltd., for supplying the ferrite tubes.

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A Technique for Minimizing Hysteresis in a 35-DB Ferrite Variable Attenuator*

A requirement arose for a low-power microwave transmitter, the output power of which could be controlled over 40 db with a reset accuracy of 0.5 db for single frequency operation.

The experimental arrangement used is shown in Fig. 1.

The electronically variable short circuit is shown in Fig. 2. This has been described by Scharfman.¹

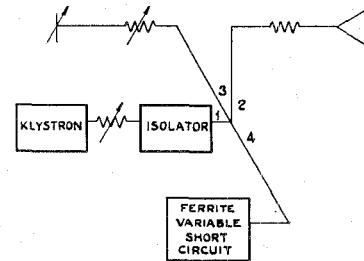


Fig. 1—Variable attenuator.

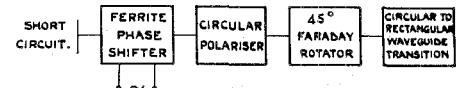


Fig. 2—Electronically variable short circuit.

For one sense of circular polarization the slope of the phase shift vs field curve becomes zero as the ferrite saturates, but for the other sense this is not so marked. This is shown in Fig. 3. Fig. 4 shows the attenuation between arms 1 and 2 vs relative phase difference between arms 3 and 4 of the magic T. It can be seen that the slope of the attenuation vs relative phase shift characteristic curve is extremely steep at the maximum attenuation point. Consider a negatively circularly polarized wave fed into a ferrite loaded section which is subjected to a field sufficiently large to saturate the ferrite. This corresponds to the point *P* in Fig. 3. The attenuator and short circuit in arm 3 can now be adjusted to give maximum attenuation between arms 1 and 2. This

* Received by the PGMTT, October 9, 1958.

¹ H. Scharfman, "Three new ferrite phase shifters," *Proc. IRE*, vol. 44, pp. 1456-1459; October, 1956.

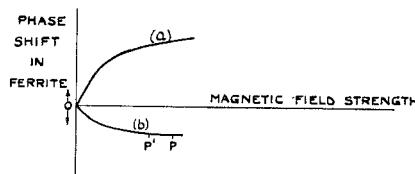


Fig. 3—Phase shift vs applied field. (a) Positively polarized wave. (b) Negatively polarized wave.

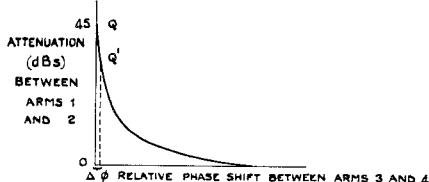


Fig. 4—Attenuation vs relative phase shift.

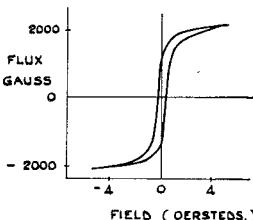


Fig. 5—Hysteresis loop of typical microwave ferrite.

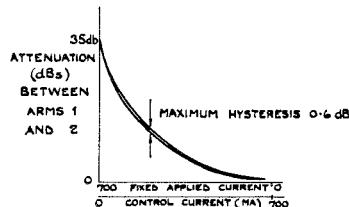


Fig. 6—Attenuator characteristic.

corresponds to the point Q in Fig. 4. If the magnetic field on the ferrite is now changed by a relatively large amount (to point P' in Fig. 3), the actual change in phase in arm 4 is quite small. This results in a small relative phase shift between arms 3 and 4 and this in turn causes the attenuation between arms 1 and 2 to change to the amount corresponding to Q' (Fig. 4). Thus it is obvious that on plotting a curve of attenuation between arms 1 and 2 vs magnetic field applied to the ferrite, a characteristic is obtained whose slope near the maximum attenuation point Q is considerably less steep than that of the curve of Fig. 4. This is shown in Fig. 6.

Since hysteresis is very small near saturation its effect near the steep part of the characteristic of Fig. 6 is very small. Below saturation the hysteresis of the ferrite is more marked (Fig. 5), but since the slope of the characteristic of Fig. 6 is much smaller when the applied field decreases, the effect of this increase in hysteresis is minimized. The final curve for the attenuator is shown in Fig. 6 where it can be seen that the maximum hysteresis measured corresponds to 0.6 dB.

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$\theta = \pi/2 + \phi$, the over-all matrix becomes

$$\begin{bmatrix} -\sin \phi & jZ_1 \cos \phi \\ j(1/Z_1) \cos \phi & -\sin \phi \end{bmatrix} \begin{bmatrix} 1 & 0 \\ j(1/Z_2) \tan \phi & 1 \end{bmatrix} \\ \cdot \begin{bmatrix} -\sin \phi & jZ_1 \cos \phi \\ j(1/Z_1) \cos \phi & -\sin \phi \end{bmatrix}$$

which when multiplied gives

$$\begin{bmatrix} \sin^2 \phi - \cos^2 \phi + (Z_1/Z_2) \sin^2 \phi \\ -jZ_1 \sin \phi \cos \phi (2 + Z_1/Z_2) \\ -j(1/Z_1) \sin \phi \cos \phi (2 - Z_1/Z_2 \tan^2 \phi) \\ \sin^2 \phi - \cos^2 \phi + (Z_1/Z_2) \sin^2 \phi \end{bmatrix}.$$

The insertion loss is given by²

$$\begin{aligned} L &= 10 \log_{10} \{ 1 + 1/4[(A - D)^2 - (B - C)^2] \} \\ &= 10 \log_{10} \{ 1 + 1/4[(2/Z_1 - 2Z_1 - Z_1^2/Z_2 \\ &\quad + 1/Z_2) \sin \phi \cos \phi - 1/Z_2 \tan \phi]^2 \} \\ &= 10 \log_{10} (1 + m^2/4) \end{aligned} \quad (1)$$

where

$$mZ_2 = R \sin \phi \cos \phi - \tan \phi \quad (2)$$

$$R = 2Z_2/Z_1 - 2Z_1Z_2 - Z_1^2 + 1. \quad (3)$$

A graph of the magnitude of $|m|Z_2$ is shown in Fig. 2.

$R = 1$ gives the maximally flat case with a zero derivative at the origin.

For R greater than 1, a triple peaked response is obtained.

Using some simple trigonometric substitutions it can be shown that $\phi_2 = 2\phi_1$; also,

$$m_1 Z_2 = \tan \phi_2 \left(\frac{1 - \cos \phi_2}{1 + \cos \phi_2} \right) \quad (4)$$

and

$$R = \frac{2}{\cos^2 \phi_2 + \cos \phi_2} \quad (5)$$

where ϕ_1 is the value of ϕ for worst reflections in the pass band, ϕ_2 is the band edge, and m_1 is the worst value of m in the pass band.

The quantity m_1 is related to the worst voltage standing wave ratio S by

$$m_1 = \frac{S - 1}{\sqrt{S}} \quad (6)$$

and the bandwidth is given by

$$BW = 2\phi_2/90. \quad (7)$$

A graph of $m_1 Z_2$ as a function of bandwidth is shown in Fig. 3.

As an example, suppose it is desired to design a stub support for a coaxial line to have a standing wave ratio of no greater than 1.05 over as wide a frequency band as possible. Because of voltage breakdown considerations it is decided that the largest value Z_2 may have is one. Then from (6), $m_1 = 0.0488$, and from Fig. 3 the bandwidth is 70.4 per cent or a frequency ratio of 2.09:1. R is determined from (5), and Z_1 from (3).

The required value of Z_1 for various values of Z_2 is plotted as a function of bandwidth in Fig. 4. This graph shows that the diameter of the quarter-wave transformers is rather critical. The desired Z_1 is only slightly smaller than the zero bandwidth case.

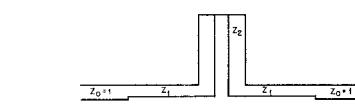


Fig. 1—Coaxial broad-band stub.

The $ABCD$ matrix of the stub plus transformers is

$$\begin{bmatrix} \cos \theta & jZ_1 \sin \theta & 1 & 0 \\ j(1/Z_1) \sin \theta & \cos \theta & -j(1/Z_2) \operatorname{ctn} \theta & 1 \end{bmatrix} \\ \cdot \begin{bmatrix} \cos \theta & jZ_1 \sin \theta \\ j(1/Z_1) \sin \theta & \cos \theta \end{bmatrix}$$

where θ is electrical length of each quarter-wave transformer and the stub. If we let

* Received by the PGMTT, October 27, 1958. The research in this document was supported jointly by the Army, Navy, and Air Force under contract with Mass. Inst. Tech.

¹ G. L. Ragan, "Microwave Transmission Circuits," M.I.T. Rad. Lab. Ser., McGraw-Hill Book Co., Inc., New York, N. Y., vol. 9, pp. 173-176; 1948.

² R. M. Fano and A. W. Lawson, "Microwave Transmission Circuits," M.I.T. Rad. Lab. Ser., McGraw-Hill Book Co., Inc., New York, N. Y., vol. 9, ch. 9 and 10; 1948.